



# DFT+U study of structural and electronic properties in $\text{La}_{0.12}\text{Sr}_{0.88}\text{TiO}_3$ : implications for solid oxide fuel cell

Estudio DFT+U de las propiedades estructurales y electrónicas del  $\text{La}_{0.12}\text{Sr}_{0.88}\text{TiO}_3$ : implicaciones para las celdas de combustible de óxido sólido

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## Abstract

The results of comparative calculations for bulk properties and electronic structure of the  $\text{SrTiO}_3$  (STO) perovskite compound and STO doped with lanthanum  $\text{La}_{0.12}\text{Sr}_{0.88}\text{TiO}_3$  (LST-12) are discussed using the ab initio density functional theory (DFT). Furthermore, a comprehensive study of the electronic structure of LST-12 with spin polarization configuration was performed. A full Bader charge analysis was employed to study the charge distributions associated with each ion. Finally, on a  $4 \times 2 \times 2$  supercell of  $\text{La}_{0.125}\text{Sr}_{0.875}\text{TiO}_{3-\delta}$  an oxygen vacancy-site was introduced to elucidate the effect on the electronic and structural properties.

**Keywords:** SOFC, anode, DFT, Bader, vacancy.

## Resumen

Este trabajo presenta un estudio comparativo de las propiedades estructurales y la estructura electrónica de la perovskita  $\text{SrTiO}_3$  (STO) y su variante dopada con lantano,  $\text{La}_{0.12}\text{Sr}_{0.88}\text{TiO}_3$  (LST-12), utilizando cálculos por primeros principios basados en la teoría del funcional de la densidad (DFT). En el compuesto LST-12, se realizó un análisis detallado de la estructura electrónica utilizando una configuración con polarización de espín. Además, se examinó la distribución de carga asociada a cada ion mediante un análisis de carga de Bader. Finalmente, se introdujo una vacancia de oxígeno en una supercelda  $4 \times 2 \times 2$  de  $\text{La}_{0.125}\text{Sr}_{0.875}\text{TiO}_{3-\delta}$  para evaluar su impacto en las propiedades del material.

**Palabras clave:** SOFC, ánodo, DFT, Bader, vacancias.

## Introduction

Solid oxide fuel cells (SOFCs) are considered as one of the most efficient and environmental-friendly technologies to address some of challenges that energy industry is currently facing. The main advantages of SOFCs include high energy conversion efficiency and big fuel flexibility. A typical SOFC single cell consists of three major components: a cathode, an anode and a ceramic electrolyte which is usually a solid oxygen ionic conductor. On the cathode side,  $\text{O}_2$  is reduced and incorporated into the electrolyte defects, namely oxygen vacancies, resulting in the formation of  $\text{O}^{2-}$  ions. This process will be driven by the difference of chemical potential, which is established between two opposite electrodes, to  $\text{O}^{2-}$  ions move through the electrolyte to the anode where these can oxidize the fuels. All the electrons generated during the electrochemical oxidation process will be released to the external circuit for power generation. The look to high-performing and stable materials for Intermediate Temperature SOFCs (IT-SOFCs) is ongoing. Nickel-zirconia cermet is commonly used as fuel electrode; this material has excellent performance when using  $\text{H}_2$  or syngas fuels but results unsuitable for direct use of hydrocarbon fuels due to the inherent carbon deposition and the irreversible sulfur poisoning. In this sense, transition metal oxides with  $\text{ABO}_3$ -type perovskite structure are currently being researched as potential IT-SOFC anodes due to their mixed ionic-electronic conductivity, thermal expansion coefficients, and tolerance to sulphur.

Strontium titanate, STO, is one such material that has been found to be a versatile compound with excellent stability in reducing atmospheres (Sharma et al., 2024). However, it has a well know band-insulating electronic structure due to the formal  $3d^0$  configuration of  $\text{Ti}^{4+}$

cations (Carballo-Córdova et al., 2019; Duan et al., 2017; Piskunov et al., 2004, 2005; Shein et al., 2005). In order to improve the electronic properties, it is possible to incorporate charge carriers through oxygen vacancies (OV) or substitution of  $\text{Sr}^{+2}$  by a trivalent ion such as  $\text{La}^{+3}$  (Choi et al., 2014).

Several experimental studies have demonstrated the suitable electronic properties of La-doped STO, which are applicable to electrochemical devices such as SOFC (Burnat, Heel, Holzer, Kata, et al., 2012; Burnat, Heel, Holzer, Otal, et al., 2012; Cheng et al., 2021; Choi et al., 2014; Deshmukh et al., 2020; Fagg et al., 2001; Hashimoto et al., 2005; Savaniu & Irvine, 2011; Yurkiv et al., 2015; S. L. Zhang et al., 2014). In fact, Presto et al. (2018) point out that  $\text{La}_{0.1}\text{Sr}_{0.9}\text{TiO}_3$  is presented as an alternative that optimizes electrochemical properties while maintaining structural stability.

By contrast, DFT calculations provide an important means of corroborating the hypotheses derived from experimental data. DFT enables the calculations of various parameters that are associated with different physical properties of the compound. These include geometric and optical properties, cell geometry, electronic band structure, the energy formation of defects and the density of states.

In the present study, first principles approach based on DFT within PBE-GGA approximation, as implemented in VASP code, was employed. Structural and electronic properties of pure  $\text{SrTiO}_3$  and La-doped  $\text{SrTiO}_3$  were investigated and compared in order to elucidate the impact of La-doping on the electrical properties of STO. In addition, a comprehensive Bader charge analysis was performed. To the best of our knowledge, this is the first report that employs a combination in procedures such as “Bader analysis” and “Small perturbation” analysis to achieve a more complete understanding of the electronic structure of stoichiometric and non-stoichiometric LST.

## Computational details

DFT calculations were performed by using Vienna Ab Initio Simulation Package (VASP) (Kresse & Hafner, 1993). The properties have been calculated by using the projector-augmented-wave method (PAW) (Blöchl, 1994) developed by Blöchl to describe the ion-electron interactions. The kinetic energy cut off  $K_{\text{max}} = 640$  eV was used for the plane-wave basis set during energy convergence. In self-consistent calculations, the total energy convergence was set to 0.01 eV. An exchange-correlation function Perdew-Burke Ernzerh (PBE) of generalized gradient approximation (GGA) (Perdew et al., 1996) was adopted for calculation of total ground state energy as well as for optimizing the structural cells internal parameters. A local Hubbard- U correction of 6.3 eV and 4 eV was also added to La and Ti atoms, respectively.

The Monkhorst-Pack's k-points grid of  $7 \times 7 \times 7$  was used in the irreducible Brillouin zone (IBZ) for electronic structure calculations.

Compounds structures were described by using different supercell sizes:  $2 \times 2 \times 2$  and  $4 \times 2 \times 2$  supercells of the primitive cell. The  $2 \times 2 \times 2$  cell was used to study the properties of STO and LST-12 whereas  $4 \times 2 \times 2$  cell was used to study the compound with oxygen vacancies ( $\text{V}_\text{O}$ ).

A series of lattice constants was set to obtain the total energy  $E$  and the corresponding primitive cell volume  $V$ . The lattice constants for the minimum total energy are the equilibrium lattice parameters, which was used to calculate the elastic constants. And then, the calculated  $E$ - $V$  data was fitted to the Birch-Murnaghan equation of state (Birch, 1947).

For all the calculations, only ionic positions were relaxed during geometry optimization

keeping the cell parameters fixed.

Charge Density Difference (CDD) was calculated by taking the difference of the charge density of the system and the isolated reference charge densities of atoms comprising the system according to equation (1). The CDD maps was plotted using VESTA software.

$$\Delta\rho = \rho_{LST} - \rho_O - \rho_{La} - \rho_{Sr} - \rho_{Ti} \quad (1)$$

Oxygen vacancy ( $O_v$ ) defects were created by removing one oxygen atom from the supercell. The OV formation energy was computed according to equation (2) (Brown et al., 2018; Ha & Alexandrova, 2016; Su & Sun, 2015; L. Zhang et al., 2016).

$$\Delta E_{f,vac} = E_{O_v} - E_{perf} - \frac{1}{2} \mu_{O_2} \quad (2)$$

where  $E_{perf}$  is the total DFT energy of the defect-free supercell,  $E_{O_v}$  is the total DFT energy of the supercell with a single defect, and  $\mu_{O_2}$  is the gas-phase chemical potential of molecular oxygen ( $\mu_{O_2} = -8.381$  eV ( $-0.616$  Ry))(Johnston et al., 2004; L. Zhang et al., 2016).

## Results and discussion

### Bulk properties

Simple cubic perovskite structure of type  $ABO_3$ , in the space group  $Pm3m$ , is the most common and stable phase of STO and LST-12 (Deshmukh et al., 2020; Hashimoto et al., 2005; Howard et al., 1989; Le et al., 2021). Figure 1 shows the 2x2x2 supercell of LST-12, in which, one Sr atom on site “A” has been replaced by La atom. The unit cell consists in 8 Ti (blue balls), 7 Sr (green balls), 1 La (brown balls) and 24 O (red balls) atoms and contains eight oxygen corner-sharing octahedra in a cubic arrangement.

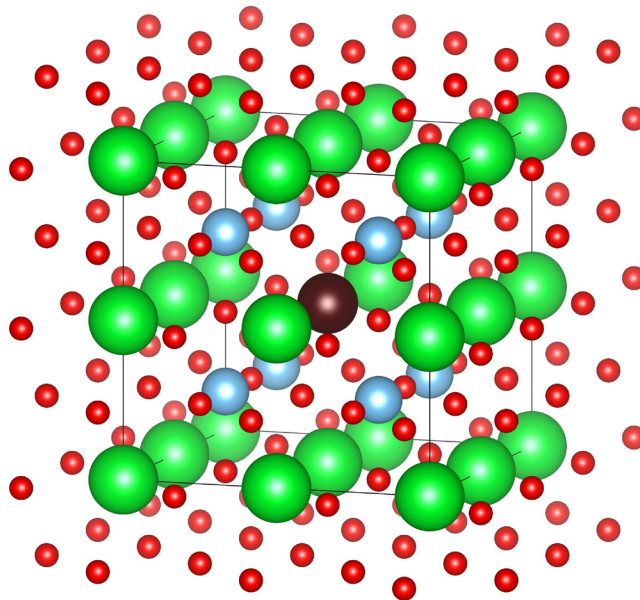


Fig. 1: 2x2x2 Supercell of the calculated structure of LST-12 (40-atom unit cell).

At the first stage, the equilibrium states for non-magnetic STO and magnetic LST-12 compounds were determined by minimizing the total energy. Table 1 lists the lattice constants of both compounds. The calculated parameters are overestimated in respect to the experimental values, but the difference does not exceed 2%.

Phase	Lattice constant		B	E <sub>0</sub>	dB/dP
	Å		GPa	eV	
	Calculated	Experimental			
LST-12	3.98	3.91	162	-301.64	4.20
STO	3.97	3.95	161	-299.21	4.21
R <sup>2</sup>	0.99997				

Table 1: Calculated and experimental values of structural properties.

A graph of energy ( $E$ ) versus volume ( $V$ ) per formula unit for STO and LST-12 is given in figure 2. By fitting the corresponding  $E$ - $V$  data to the Birch-Murnaghan equations of state it was obtained the isothermal bulk modulus  $B_0$  and its pressure derivative  $dB/dP$ . The obtained values are listed in table 1. As observed, the elastic properties are remained practically unchanged in the phase doped with La. All obtained values are in good agreement with the experimental data (Adewale et al., 2020).

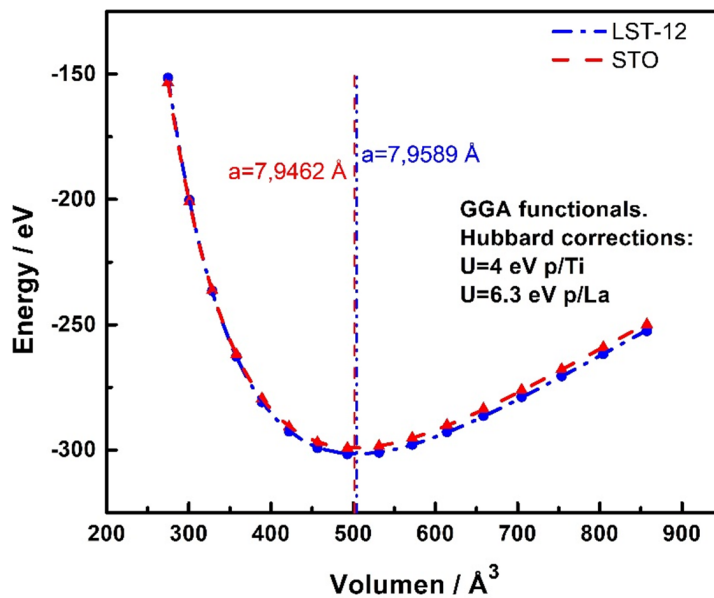


Fig. 2: Relations of dependence of the total energy on the lattice parameter for the magnetic phases of LST-12 and STO perovskites.

## Electronic properties

Figure 3 shows the comparison of calculated densities of states (DOS) for STO and LST-12. In the ground state, STO is an insulator with a direct band-gap of 2.27 eV. This value is underestimated with respect to the reported experimental band-gap (3.75 eV) (Van Benthem et al., 2001). The LST-12 compound displays semi-metallic behaviour as the Fermi energy ( $E_F$ )

level lies near the bottom in the conduction band (CB). This means the material is an electron charge carrier (n-type).

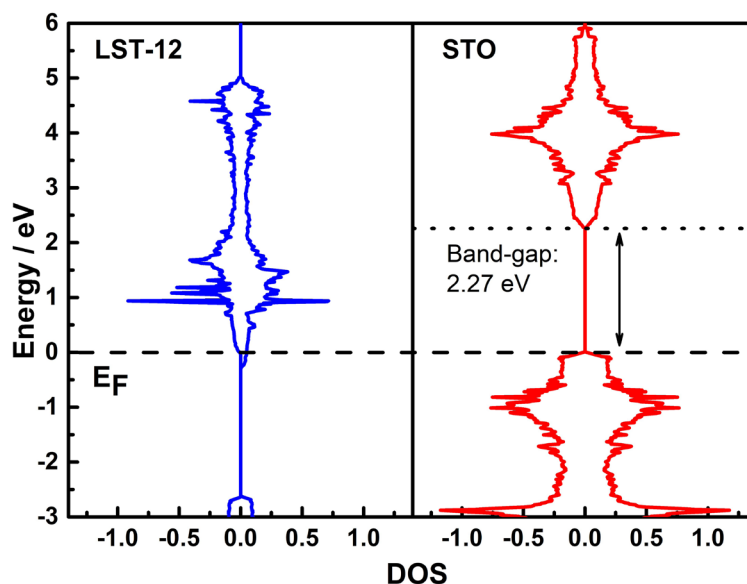


Fig. 3: The calculated total DOS for LST-12 and STO perovskites.

The contributions of each atom of LST-12 to the band structure lines cluster were analysed by computing the projected partial density of states. The total DOS and the band structure are presented in figure 4 (a) while the projected DOS are shown in figure 4 (b). The band structure has been plotted along the high symmetry points  $\Gamma$ -X-M- $\Gamma$ -R-X-M of the Brillouin zone.

As these are displayed in figure 4 (b), the low-energy region around the gap is constituted by a valence band (VB) of mostly O-2p orbitals, while the conduction band (CB) is primarily characterized by Ti-3d orbitals.

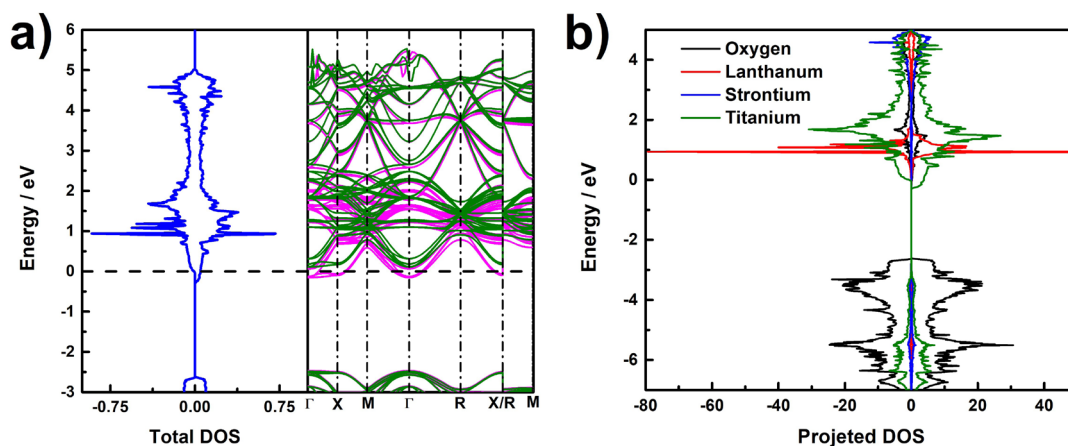


Fig. 4: Calculated total DOS with band structure (a) and (b) projected DOS for LST-12 perovskite.

A magnetic behaviour was confirmed with different spin-up and spin-downs of DOS orbitals. Accordingly, the total magnetic moment calculated was  $0.792 \mu_B$  and over each atom results  $0.001 \mu_B$  for Sr,  $0.023 \mu_B$  for La,  $0.124 \mu_B$  for Ti and  $-0.01 \mu_B$  for the oxygen.

Ti-ions in both STO and LST-12 present an octahedral coordination with six O-ligands surrounding them. This is confirmed in figure 5 in which the crystal field splitting is observed. The electrons in d-orbitals experience repulsion with the ligand electrons. This repulsion is stronger in the case of  $dx^2-y^2$  and  $dz^2$  orbitals as these point along the axes towards the direction of the ligand. Hence, they present higher energy than the average in the spherical crystal field. On the other hand,  $dxy$ ,  $dyz$  and  $dxz$  orbitals experiment lower repulsions as they lie along the diagonals of the three planes. Therefore, these three orbitals have less energy than the average energy in the spherical crystal field. Thus, the 5-fold degenerate d-orbitals are divided into two energy levels:  $t_{2g}$ -set of three orbitals ( $dxy$ ,  $dyz$  and  $dxz$ ) with low energy and  $e_g$ -set of two orbitals ( $dx^2-y^2$  and  $dz^2$ ) with high energy.

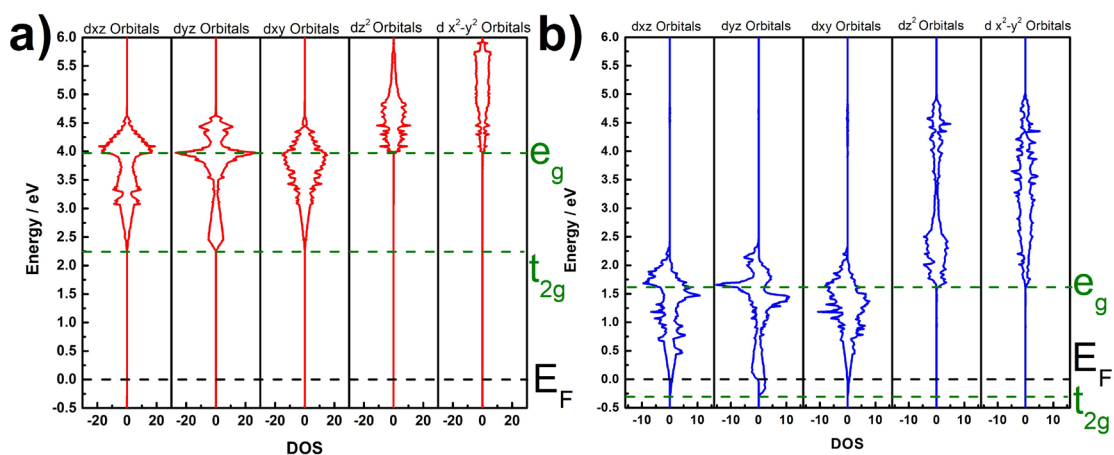


Fig. 5: The calculated DOS of Ti 3d states and their decomposition into the components  $t_{2g}$  and  $e_g$  for the (a) STO and (b) LST-12 perovskites.

Both compounds present different crystal splitting behaviour with different  $10 Dq$ . Figure 5 (a) shows the projected DOS of Ti-3d in STO and figure 5 (b) shows the projected DOS of Ti-3d in LST-12. As it can be seen, the values of crystal field splitting  $10 Dq$  are 1.73 eV and 1.9 eV in STO and LST-12, respectively. These values are underestimated respect the reported experimental values (Fan et al., 2019). La-doping in the A-site causes a slight increase in the splitting of the d orbitals, being more unfavourable to put electrons in higher energy sets of orbitals instead of pairing them up.

Bader charge analysis (Henkelman et al., 2006; Sanville et al., 2007; Tang et al., 2009) was used to analyse the charge density of the LST-12 phase and calculate the associated charge to each ion. This analysis allows to describe quantitatively the ionic/covalent character of the compound since under Bader's scheme the electron topology is characterized in terms of its critical points (CPs) which are defined as the zeros of the gradient of the electron density.

Atom	$Q_B$ (e)	Nominal ionic valences	VB ( $\text{\AA}^3$ )	$\delta$ ( $\text{e}/\text{\AA}^3$ )
Oxygen	-1.268 (Min)	-2	13.0694	0.097
	-1.1954 (Max)		12.6558	0.094
Strontium	1.5866 (Min)	+2	16.3651	0.097
	1.5979 (Max)		16.4074	0.097
Lanthanum	2.1922	+3	18.3964	0.119
Titanium	2.2078	+4	7.7331	0.284

Table 2. Parameters resulting from Bader analysis.

The Bader charge of Ti and O differs significantly from the formal ionic charges of  $\text{ABO}_3$  perovskite:  $\text{Ti}^{4+}$  and  $\text{O}^{2-}$ . This can be explained by the significant overlap between the Ti 3d and O 2p orbitals, which results in a partial covalent O-Ti chemical bonding. In contrast, there is practically no bonding of O with Sr and La atoms. The charges of Sr and La remain close to the formal +2e and +3e respectively.

Figure 6 shows the CDD map on planes of interest. The positive CDD implies charge gain, while a negative CDD implies loss of electric charge. The colour reference ranges from blue in the negative extreme to red in the positive extreme.

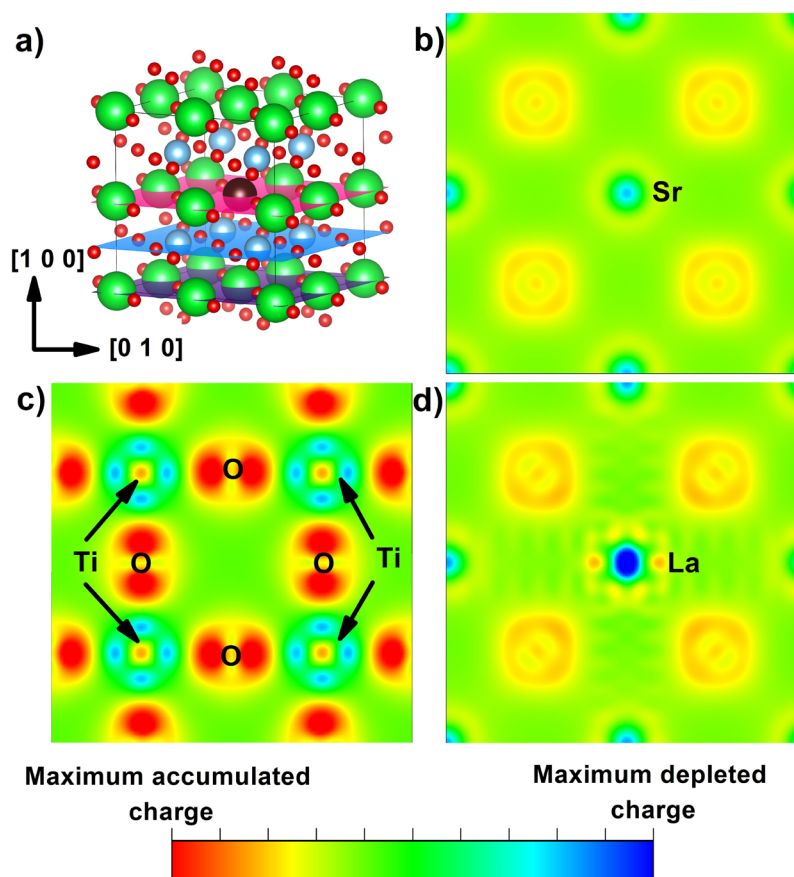


Fig. 6: Difference electron density plots for different cross sections; (a) structural unit of the cubic LST-12 perovskite crystal cross-sectioned by (1 0 0) plane, (b) CDD maps for Sr-O (1 0 0) cross sections, (c) CDD maps for Ti-O (1 0 0) cross sections and (d) CDD maps for La-O (1 0 0) cross sections

Analysis of the CDD maps confirms the Ti-O covalent bonding effect mentioned above. Figure 6 (c) shows the concentration of the electronic density between Ti and O ions. In contrast, the CDD maps drawn for the AO-(1 0 0) cross sections (Figure 6 (b) and (d)) show no trace of the covalent bonding between the oxygen atom and Sr or La. The A cations on the AO-terminated surface demonstrate considerable polarization. Slight differences can be observed between the planes containing La and Sr. Figure 6 (b) shows that the electronic charge is spherically distributed around the sites occupied by O atoms, which is consistent with a perfect ionic bond. In contrast, in Figure 6 (d) it shows an asymmetric charge accumulation around oxygen and a greater depleted charge in the La position, due to the doping with these atoms.

When an extra electron is added to the material, in this case from the La donor, it will reside in two possible configurations. If the symmetry of the lattice is unbroken, the extra electron will be delocalized at the CB minimum. However, if the extra electron couples to the ions in the crystal a local distortion of the lattice is induced. This distortion produces an electric field that acts on the electron and generate a feedback mechanism that alters the energetics and dynamics of the charge carrier and may induces self-trapping. In consequence, a polaron is formed. The strong interaction of virtual phonons and electrons leads to problems in charge transport.

To study the possible formation of a quasiparticle, an intentional lattice deformation was introduced (Density Functional Perturbation Theory-DFPT) by changing the bond lengths of Ti-O (2.1 Å) and then performing a self-consistent calculation allowing the relaxation of ionic positions. The titanium charge values obtained from the Bader analysis, after the “small perturbation”, indicate that the electron behaves as a free particle in a Fermi gas. The Bader charge value results around +2.21e for all Ti positions.

## Oxygen vacancies

The oxygen vacancy formation energy in LST-12, calculated in our DFT study, yields 6.45 eV. This value is in good agreement with those reported for the undoped STO phase (Brown et al., 2018; L. Zhang et al., 2016). As can be seen in figure 7, the removal of one oxygen atom diminishes the symmetry of the crystal, leading to significant distortions in the crystal lattice due the rotation and deformation of  $\text{TiO}_6$  octahedra. The Ti-O<sub>v</sub>-Ti distance is increased to 4.02 Å while the Ti-O-Ti distances are slightly modified with respect to the distance in stoichiometric LST-12.

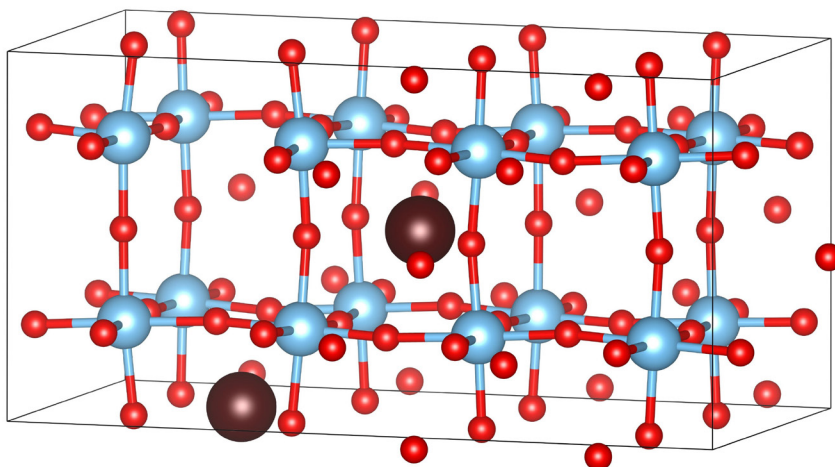


Fig. 7: Supercell 4x2x2 of the calculated structure of LST with one oxygen vacancy (79-atom unit cell)

A similar behaviour has been observed by other authors (Souto-Casares et al., 2021; Winczewski et al., 2022; L. Zhang et al., 2016) in their calculation performed for STO phase.

Figure 8 shows the comparison between the total DOS and band structure of the stoichiometric LST-12 (Figure 9 a)) and defective  $\text{La}_{0.12}\text{Sr}_{0.88}\text{TiO}_{3-\delta}$  (Figure 9 b)). As it can be exhibited, the main difference between the two calculations is the introduction of an electronic band at -0.6 eV below the Fermi level. The two electrons provided by the oxygen vacancy, mainly occupy the bands corresponding to the  $\text{d}_{z^2}$  orbitals of titanium (Figure 9 c)).

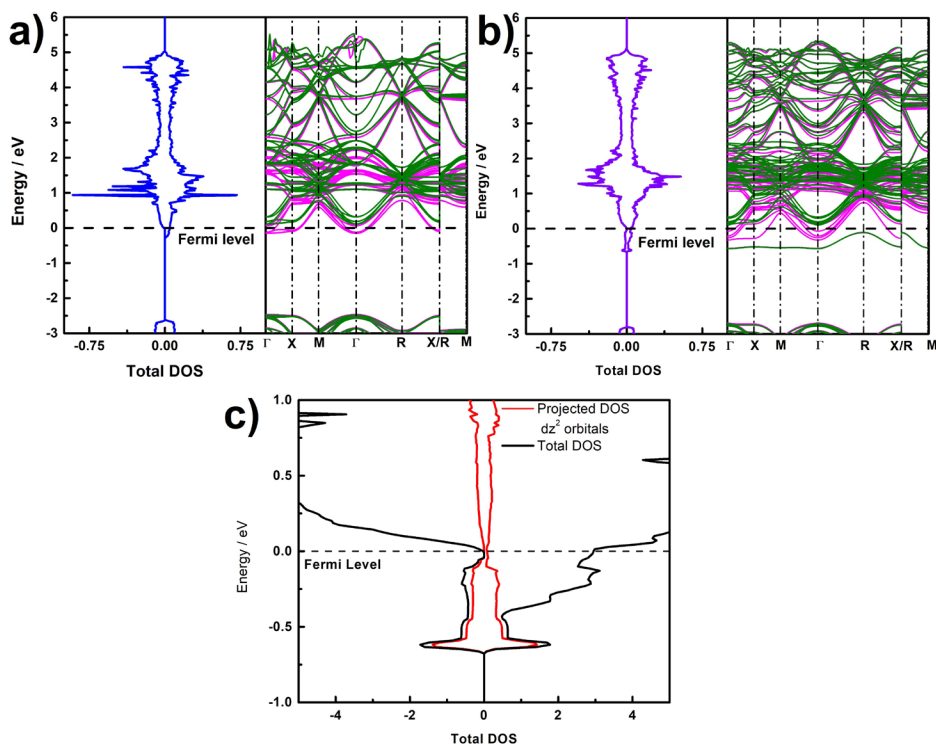


Fig. 8: Calculated DOS and band structure for (a) stoichiometric LST and (b) defective  $\text{La}_{0.125}\text{Sr}_{0.875}\text{TiO}_{2.9375}$  and (c) calculated total DOS and projected DOS of Ti  $3\text{d}_{z^2}$  states.

The Bader analysis for the system with defects shows a decrease in the charge value of the Ti near to the vacancy site to +1.83e ( $\text{Ti}_4$  and  $\text{Ti}_7$  in figure 8) while the other Ti have a charge value around +2.2e which could be attributed to localized electrons forming  $\text{Ti}^{+3}$  polarons trapped around isolated oxygen vacancies.

## Conclusions

First-principles calculations and Bader analysis were performed to examine the bulk and electronic properties of STO, LST and non-stoichiometric LST perovskite, as potential SOFC anodes. Calculations were performed with the generalized gradient approximation (GGA) taking into account the on-site Coulomb repulsive interaction (GGA + U). The results reveals that the calculated lattice constants of both compounds are in good agreement with experimental data and in previous reports available.

Regarding the electronic properties, it was demonstrated that when STO is doped with a donor such as La, it exhibits metallic-like behavior. In addition, LST displays magnetic properties with a total magnetic moment of  $\sim 0.8 \mu_B$ , whereas STO presents non-magnetic behavior. Also, it was possible to observe the crystal field splitting by Ti-projected DOS. The 10 Dq yielded 1.73 eV and 1.9 eV for STO and LST, respectively. Bader's analysis facilitated an understanding of the nature of the chemical bonds in the compound, revealing a partially Ti-O covalent bond and a strong ionic character in Sr and La, which was confirmed by CDD maps. Additionally, polaronic modes of conduction were discarded by means of an intentional lattice deformation and subsequent ionic relaxation.

Finally, the formation energy of a vacancy was calculated, and the distortion of the crystalline structure was verified. A lower symmetry was observed due to the rotation of the octahedra  $TiO_6$ . A notable decrease in the Bader charge value was evident in the Ti region close to the vacancy, which could be attributed to the formation of  $Ti^{+3}$  polarons.

The main contribution of this work was relating the electronic and structural properties at the atomic level to the expected performance as an SOFC anode. It was observed that doping with La provides the material with adequate characteristics for hydrogen catalysis and good electronic conductivity properties.

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1-Administración del proyecto, 2-Adquisición de fondos, 3-Análisis formal, 4-Conceptualización, 5-Curaduría de datos, 6-Escritura - revisión y edición, 7-Investigación, 8-Metodología, 9-Recurso, 10-Redacción - borrador original, 11-Software, 12-Supervisión, 13-Validación, 14-Visualización.